

# Gross-Neveu model as a laboratory for fermion discretization

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Lattice 2006, Tucson AZ

- Models, scheme, simulation
- Example results, Wilson:  $N = 2, 8$
- Staggered (no results yet), Outlook

# The models

Dirac  $\leftrightarrow$  Majorana fermions, free Euclidean theory:

$$\mathcal{L}_0 = \sum_{k=1}^{N/2} \bar{\psi}_k (\not{\partial} + m) \psi_k = \frac{1}{2} \sum_{k=1}^N \xi_k^\top \underbrace{\mathcal{C}(\not{\partial} + m)}_{\text{antisymmetric}} \xi_k$$

with

$$\begin{pmatrix} \psi_1 \\ \bar{\psi}_1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \xi_1 + i\xi_2 \\ (\xi_1^\top - i\xi_2^\top)\mathcal{C} \end{pmatrix}, \text{ etc, } \quad \mathcal{C}\gamma_\mu\mathcal{C}^{-1} = -\gamma_\mu^\top \quad [\Rightarrow \mathcal{C}^\top = -\mathcal{C}]$$

$\Rightarrow$  large internal ‘flavor’ symmetry  $\text{O}(N) \supset \text{U}(N/2)$

by insisting on  $O(N)$  invariance:

- only one 4-fermion interaction possible

$$\mathcal{L}_I = -\frac{g^2}{8} (\xi_i^\top \mathcal{C} \xi_i)^2$$

- renormalizable in  $D = 2$ 
  - $N = 2$ : Thirring model
  - $N \geq 3$ : original asymptotically free Gross-Neveu models
- also odd  $N$  possible
- discrete chiral invariance  $\xi_i \rightarrow \gamma_5 \xi_i$  iff  $m = 0$  (Wilson:  $\rightarrow m_c$ )
- GN with chiral  $U(1)$  needs three couplings (mixing under renormalization!)  $\Rightarrow$  disfavored for precision tests

# Finite size renormalization scheme

- $T \times L$  torus,  $\xi(x \pm T\hat{0}) = -\xi(x)$ ,  $\xi(x \pm L\hat{1}) = \xi(x)$
- IR-regularized by smallest momentum  $p_* = (\pi/T, 0)$
- formulate nonperturbative chiral renormalization conditions  
at scale  $T$
- momentum space: normalize at momentum  $p_* \rightarrow$  not discussed  
here

- position space,  $T = L$ , **PT** refers to **Wilson** fermions [ $r = 1$ ]

$$\xi_R = Z_\xi^{-1/2} \xi, \quad \check{\xi}(x_0) = a \sum_{x_1} \xi(x)$$

$$\left\langle \check{\xi}_1^\top(T/4) \mathcal{C} \check{\xi}_1(0) \right\rangle = 0 \Rightarrow m = m_c \text{ (} \rightarrow \text{tune)}$$

PT:  $am_c = -(N-1)Kg^2 + O(a^4g^2, g^4), \quad K = 0.384900179460$

$$Z_\xi = -\frac{1}{L} \left\langle \check{\xi}_1^\top(T/2) \mathcal{C} \gamma_0 \check{\xi}_1(0) \right\rangle [= 1 + O(g^4) \text{ in PT}]$$

$$g_R^2 = \frac{4}{TL} \left\langle \check{\xi}_{R,1}^\top(T/2) \mathcal{C} \check{\xi}_{R,1}(0) \check{\xi}_{R,2}^\top(T/2) \mathcal{C} \check{\xi}_{R,2}(0) \right\rangle$$

PT:  $g_R^2 = \frac{T}{T-2a} \left( g^2 + \left[ \frac{N-2}{2\pi} \ln(L/a) + c_0 + O(a) \right] g^4 + O(g^6) \right)$

$$c_0 = 0.309652\dots - (N-2) \times 0.483524\dots$$

# Simulation

$$e^{\frac{g^2}{8} a^2 \sum_x (\xi_i^\top C \xi_i)^2} = \int \prod_x d\mu(\sigma) e^{-\frac{g}{2} a^2 \sum_x \sigma \xi_i^\top C \xi_i}$$

valid if  $\int d\mu(\sigma) \begin{Bmatrix} \sigma^{2n} \\ \sigma^{2n+1} \end{Bmatrix} = \begin{Bmatrix} \frac{2n!}{n!} (2a^2)^{-n} \\ 0 \end{Bmatrix}$  for  $n \leq N/2$

standard:  $d\mu = \text{Gaussian}$  (for all  $N$ )  $\Rightarrow$

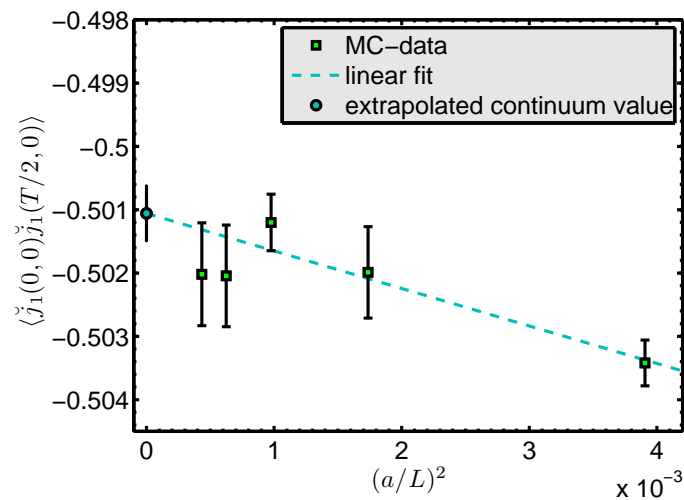
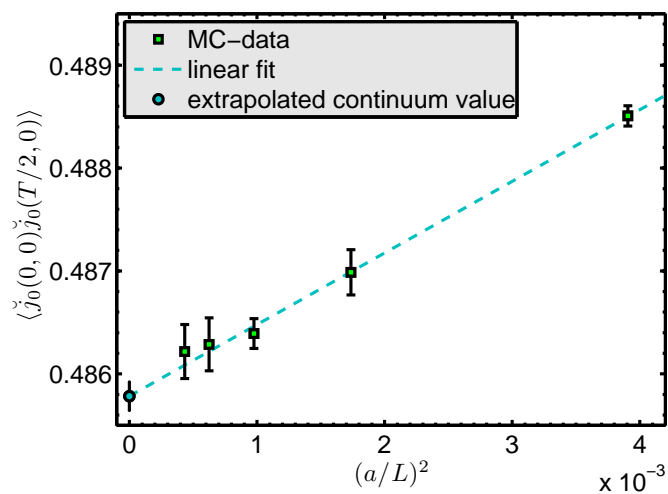
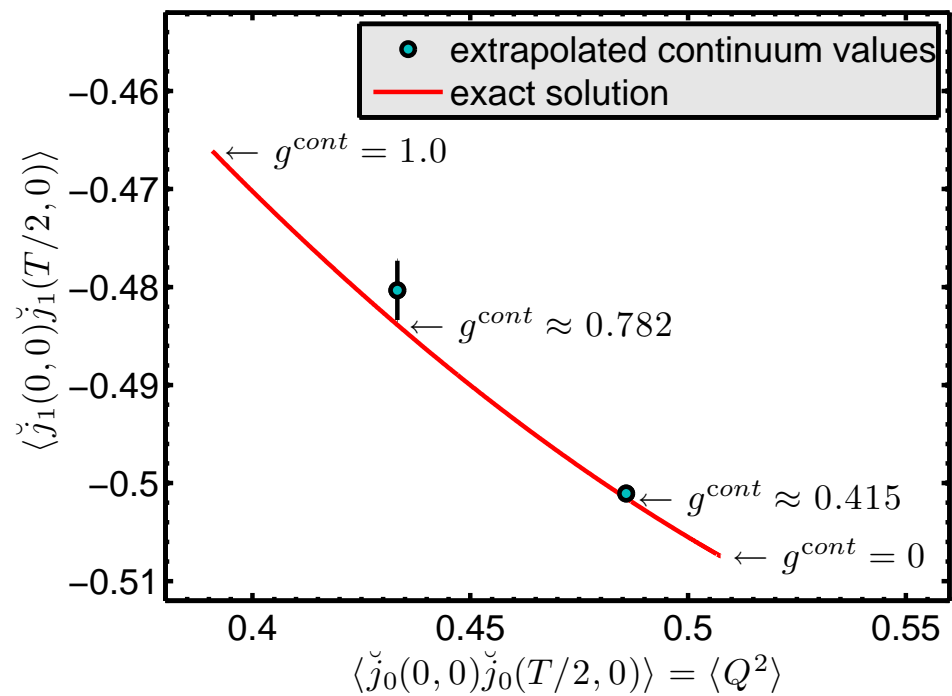
$$Z = \int \prod_x d\sigma e^{-\frac{1}{2} a^2 \sum_x \sigma^2} \underbrace{[\text{Pf } A(\sigma)]^N}_{=(\det A)^{N/2}}$$

$$A(\sigma) = \mathcal{C}(\gamma_\mu \partial_\mu + m + g\sigma - a\partial^* \partial)$$

- $A$  antisymmetric matrix, real with Majorana  $\gamma_\mu$
- for even  $N \Rightarrow 0 \leq (\text{Pf } A)^N = (\det A)^{N/2} = (\det A^\top A)^{N/4}$
- simulable by standard methods ( $N/2$  real pseudofermions)

## $N = 2$ : Massless Thirring model

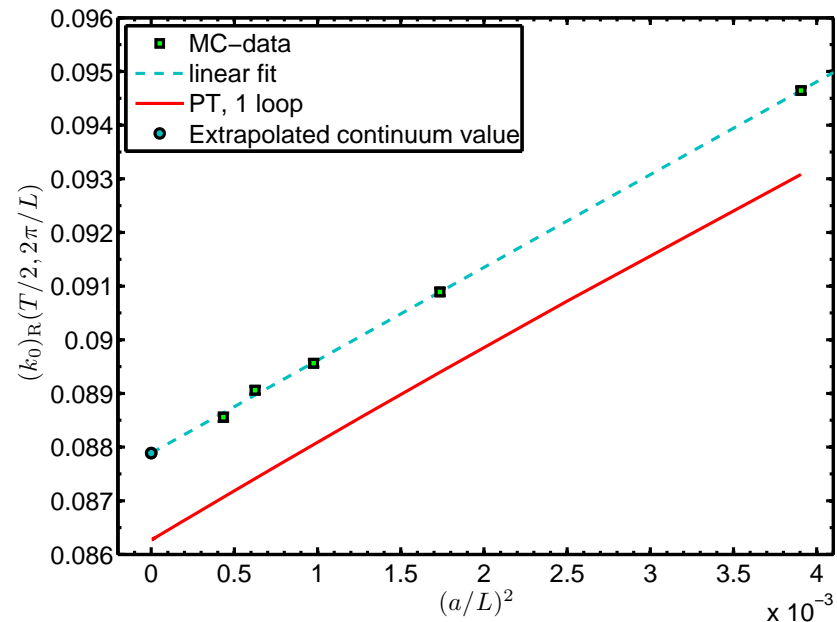
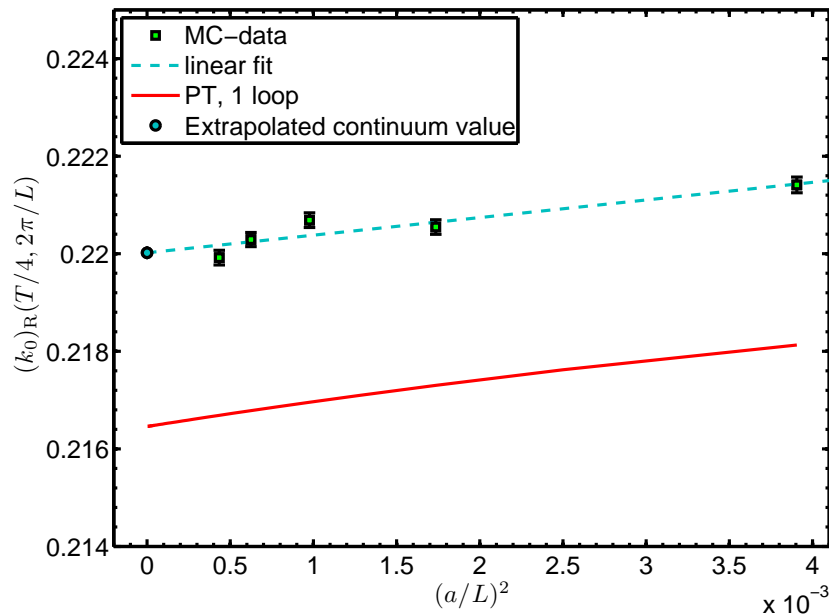
- exact solution in the continuum ( $T = L = \infty$ )
- only finite coupling renormalization, no running,  $\beta \propto (N - 2)$
- key formula:  $\det(\not{D} + i\not{A}) = \exp\left(-\frac{1}{2\pi} \int d^2x A_\mu [\delta_{\mu\nu} - \frac{\partial_\mu \partial_\nu}{\partial^2}] A_\nu\right)$
- extended to  $T, L < \infty$  (complicated by zeromodes)
- correlations of  $\check{j}_\mu(x_0, p) = \int_0^L dx_1 e^{-ipx_1} \bar{\psi} \gamma_\mu \psi$  worked out
- $\check{j}_0(x_0, 0) = Q$  charge,
- $\Rightarrow$  check with Wilson fermions
  - continuum limit at fixed bare  $g = 0.4, 0.7$   
 $L/a = 16 \dots 48$  always tuning  $m = m_c(g)$
  - use Noether current ( $Z_V = 1$ )



# N=8 Gross Neveu

- many universal correlations definable
- one example (at  $g_R = 0.38$ )

$$(k_\mu)_R(x_0, p) = -\frac{1}{Z_\xi} a \sum_{x_1} e^{-ipx_1} \langle \xi^\top(0) \mathcal{C} \gamma_\mu \xi(x) \rangle$$



# Staggered Fermions

$D$  dimensions

- naive fermions:  $2^D$  ‘tastes’ (Dirac and Majorana)
- reduction factor by “spin diagonalization”:  
Dirac:  $2^{D/2}$ , Majorana:  $2^{D/2-1}$  (none in  $D = 2$ )
- reason:  $\mathcal{C}$  not transformed away (has  $2 \times 2$  blocks)
- $D = 2$ : 2 Dirac- or 4 Majorana-tastes (  $\times 2$  to simulate)
- $m_c = 0$ , no additive renormalization
- naive 4-fermion term:

$$\frac{1}{(TL)^4} \sum_{p_1, \dots, p_4} \delta^2\left(\sum p_i\right) \underbrace{\xi_i^\top(p_1) \mathcal{C} \xi_i(p_2)}_{\text{taste-mixing!}} \xi_j^\top(p_3) \mathcal{C} \xi_j(p_4)$$

→ extra factor:  $\prod_\mu \cos(a(p_1 + p_2)_\mu/2) \cos(a(p_3 + p_4)_\mu/2)$

locally square in position space:

$$\frac{1}{8} \left[ \xi^\top \mathcal{C} \xi (x) + \xi^\top \mathcal{C} \xi (x + a\hat{0}) + \xi^\top \mathcal{C} \xi (x + a\hat{1}) + \xi^\top \mathcal{C} \xi (x + a\hat{0} + a\hat{1}) \right]$$

- $O(4N)$  Gross-Neveu in the continuum limit??
- compare universal correlations

## Summary, Outlook

- Finite size scheme for GN
- relatively easy to tune chiral (similar to Schrödinger functional in QCD), also for Wilson, + continuum limit
- easy to define universal correlations (ratios...)
- [precise universality check](#), staggered, Wilson, ... [possible?](#)
- Nonstandard fermion [algorithms](#)? Dimer? World line?  
Other factorizations of  $(\xi^\top \mathcal{C} \xi)^2$  ?