

Progress towards the phase diagram of QCD

Philippe de Forcrand

ETH Zürich & CERN

with

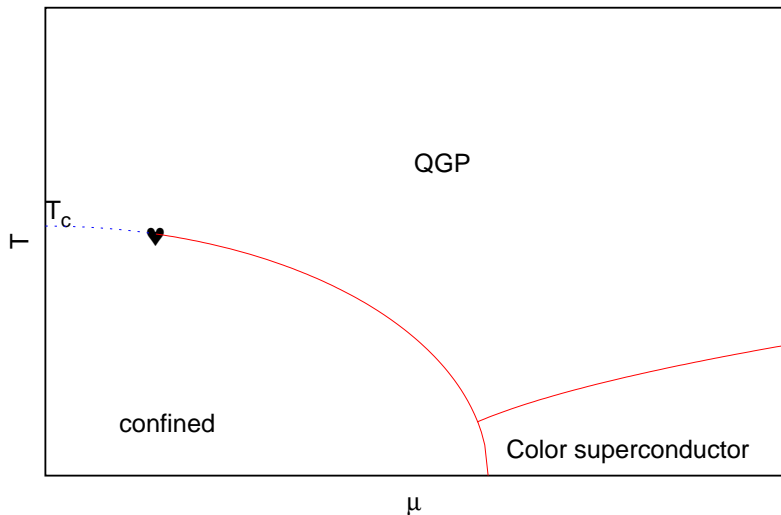
Owe Philipsen

Münster

LAT06, Tucson, July 2006

hep-lat/0607017

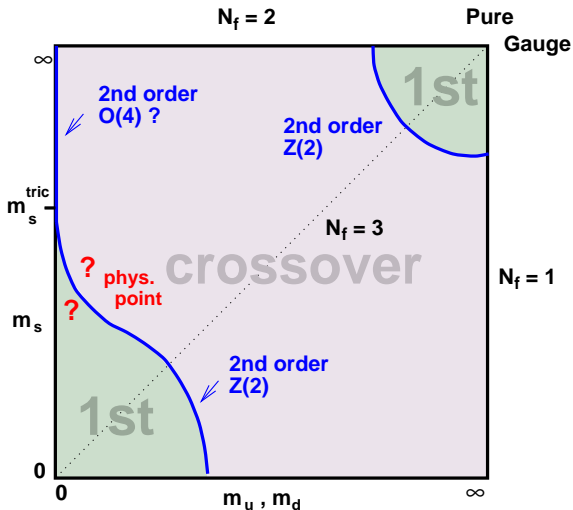
Phase diagram of QCD



Learn about QCD by generalizing to arbitrary $(m_{u,d}, m_s)$ quark masses

Phase diagram vs $(m_{u,d}, m_s)$, T and μ

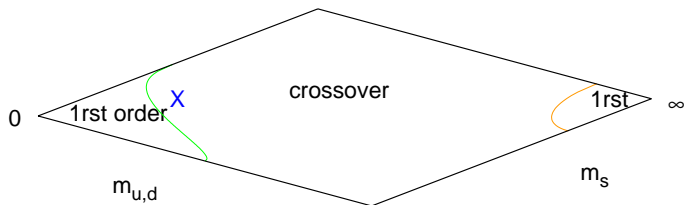
$$\mu = 0$$



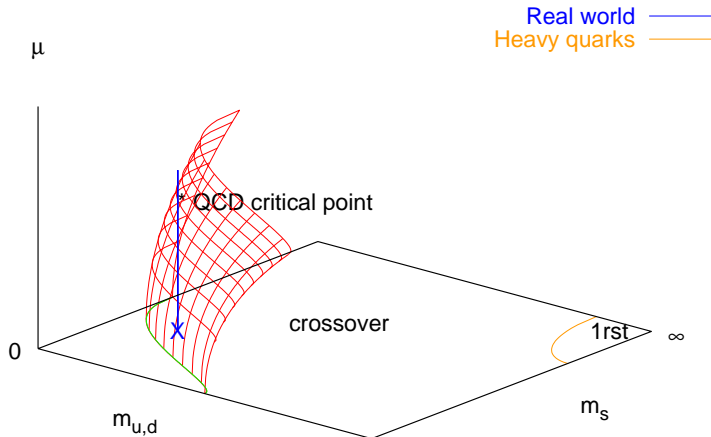
Phase diagram vs $(m_{u,d}, m_s)$, T and μ

$$\mu = 0$$

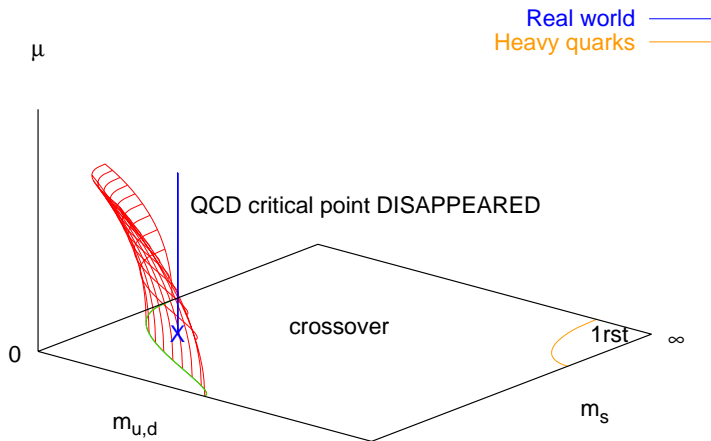
Real world ————
 Heavy quarks ————



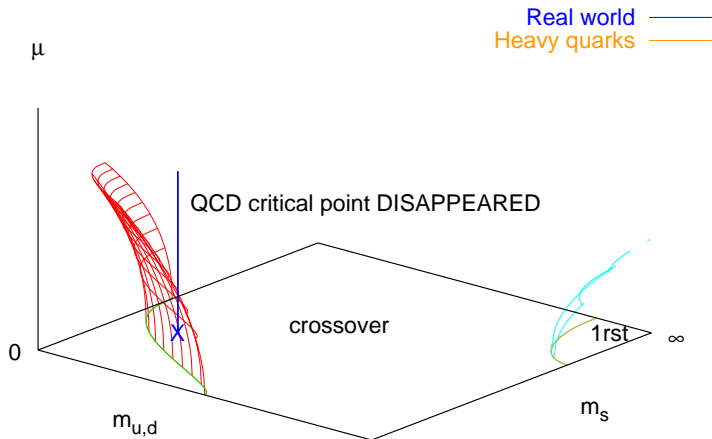
Now turn on μ

Phase diagram vs $(m_{u,d}, m_s), T$ and μ $\mu \neq 0$ 

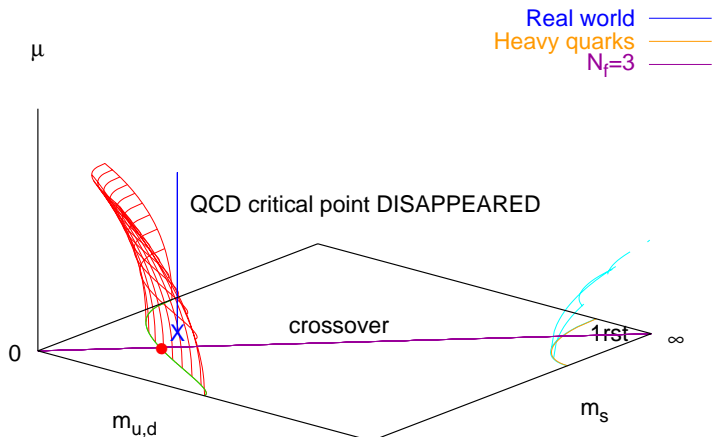
Conventional wisdom: first-order region **expands** with real $|\mu|$

Phase diagram vs $(m_{u,d}, m_s)$, T and μ 

Exotic scenario: first-order region **shrinks** with real $|\mu|$

Phase diagram vs $(m_{u,d}, m_s), T$ and μ 

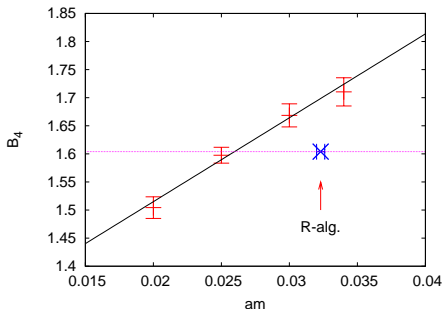
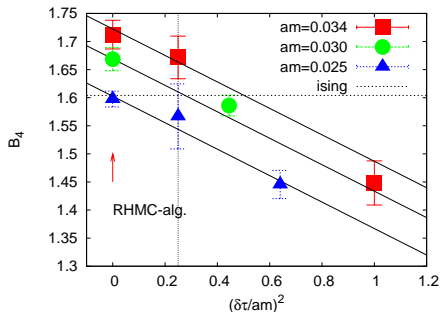
For heavy quarks, first-order region shrinks (PdF, Kim & Takaishi)

Phase diagram vs $(m_{u,d}, m_s)$, T and μ 

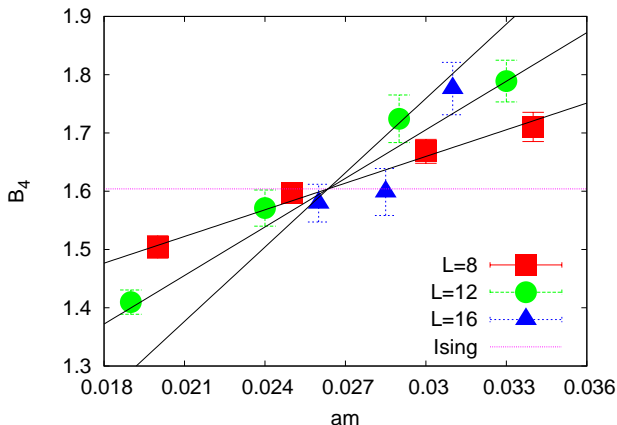
Detect 2nd order with Binder cumulant $B_4 \equiv \frac{\langle (\delta\bar{\psi}\psi)^4 \rangle}{(\langle (\delta\bar{\psi}\psi)^2 \rangle)^2}$. Check $N_f = 3$ first.

Eliminating the stepsize error: R \rightarrow RHMC

- Old R-alg. requires $\delta\tau \rightarrow 0$ extrapolation
- $\delta\tau \neq 0 \rightarrow$ bias transition towards 1st-order (Kogut & Sinclair)

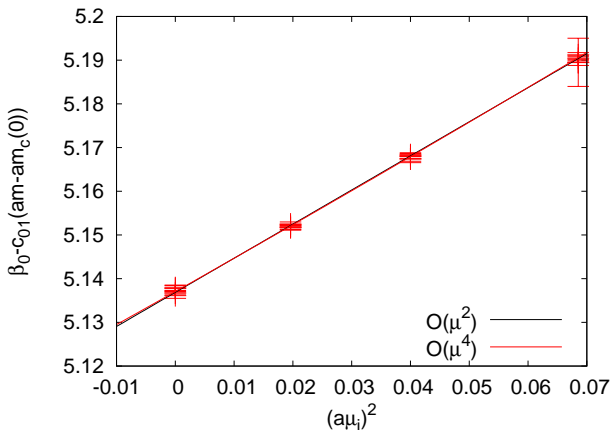


- New RHMC (Clark & Kennedy) exact and more efficient
- $N_f = 3$: critical quark mass for 2nd order P.T. at $\mu = 0$ down by $\sim 25\%$!
- Lattice renorm.? No: corresponding $\frac{m_\pi}{T_c}$ down by $\sim 10\%$

$N_f = 3, \mu = 0$: finite size scaling

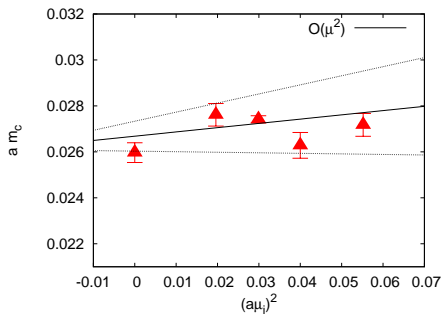
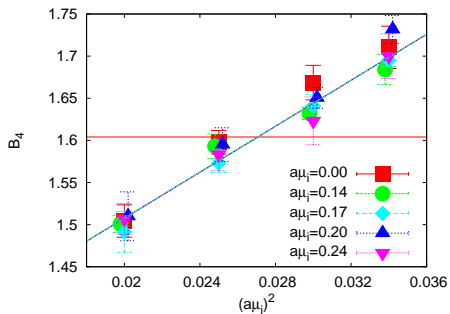
$$\xi \propto |m - m_c|^{-\nu} \rightarrow B_4(m, L) = B_4^{Ising} + bL^{1/\nu}(m - m_c)$$

Fit: $\nu = 0.67(13)$ versus $\nu^{Ising} \approx 0.63$

$N_f = 3, \mu \neq 0$: pseudo-critical temperature

Correction $\sim 10\%$ from R-algorithm:

$$\frac{T_c(\mu, m)}{T_c(0, m_0^c)} = 1 + 2.111(17) \left(\frac{m - m_0^c}{\pi T_c} \right) - 0.667(6) \left(\frac{\mu}{\pi T_c} \right)^2 + 0.23(9) \left(\frac{\mu}{\pi T_c} \right)^4 + \dots$$

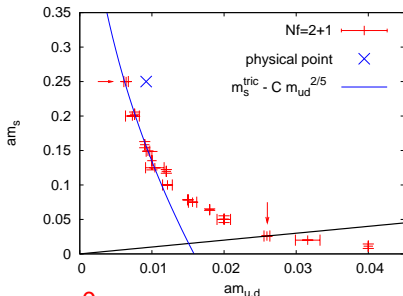
$N_f = 3, \mu \neq 0$: critical mass versus μ 

No significant variation in $a m_c \propto \frac{m_c}{T_c}$ versus $a\mu \propto \frac{\mu}{T}$

$$\Rightarrow \frac{d m_c}{d(\mu/T)^2} \approx \frac{d T_c}{d(\mu/T)^2}$$

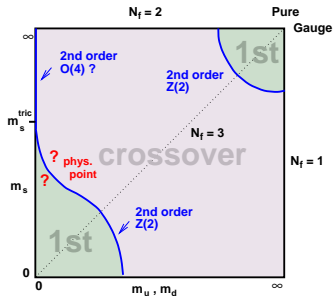
$$\frac{m_c(\mu)}{m_c(\mu=0)} = 1 - 0.7(4) \left(\frac{\mu}{\pi T}\right)^2 + \dots$$

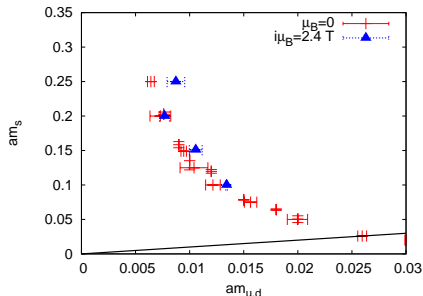
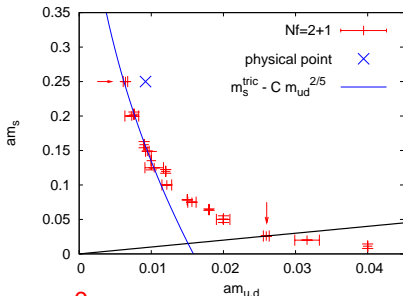
The first-order region shrinks when μ is turned on

Extend to $N_f = 2 + 1$ Line of second-order phase transitions in the quark mass plane ($m_{u,d}, m_s$) $\mu = 0$:

- data consistent with **tricritical point** at $m_{u,d} = 0$, $m_s \sim 2.8T_c$
- physical point **in crossover region**

Fodor & Katz



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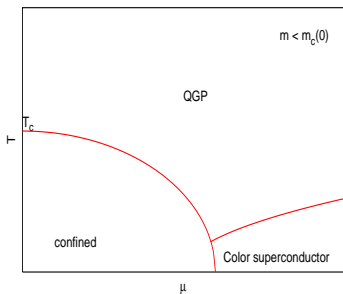
 $\mu \neq 0$ imaginary:

- am_c increases slightly \Rightarrow decreases slightly with **real μ**
- $$\Rightarrow \frac{d m_c}{d(\mu/T)^2} \lesssim \frac{d T_c}{d(\mu/T)^2}$$

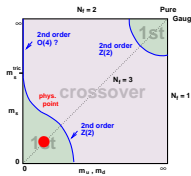
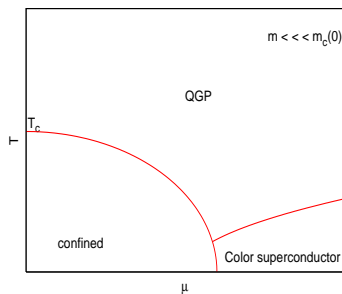
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Resulting phase diagram (simplest possibility)

Standard scenario

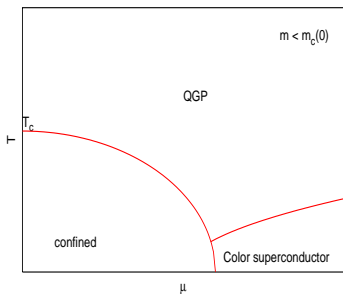


Exotic scenario

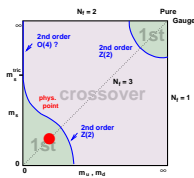
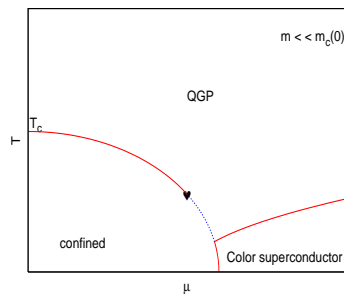


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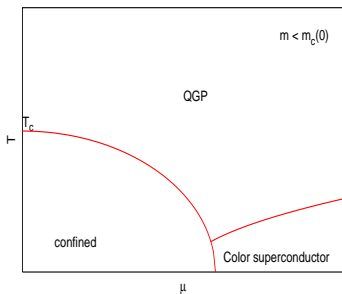


Exotic scenario

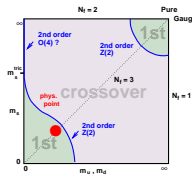
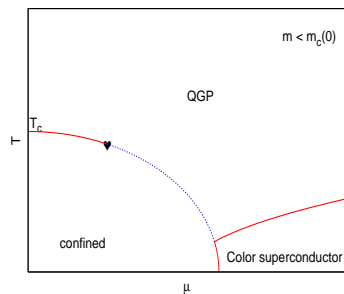


Resulting phase diagram (simplest possibility)

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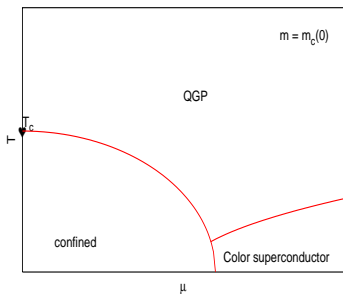


Exotic scenario

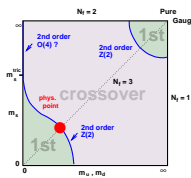
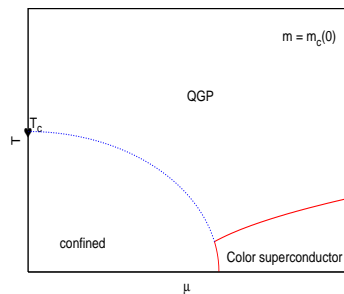


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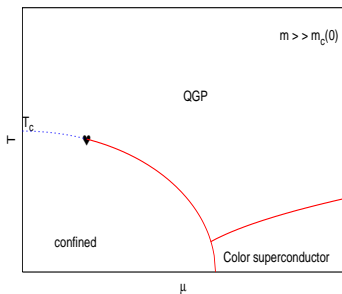


Exotic scenario

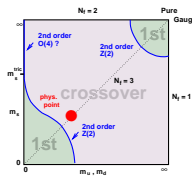
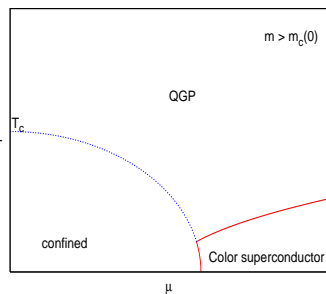


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Standard scenario



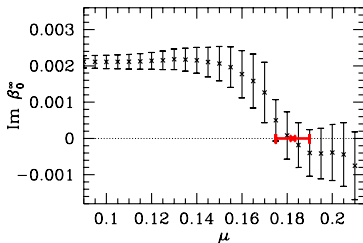
Exotic scenario



Contradiction with other lattice studies?

Caveat: at $N_f = 3$, $\mu = 0$ critical point, m_π varies by **factor** ~ 4
with Dirac discretization **Bielefeld, MILC**

- **Fodor & Katz:** $(T_E, \mu_E) = (162(2), 120(13))$ MeV ?

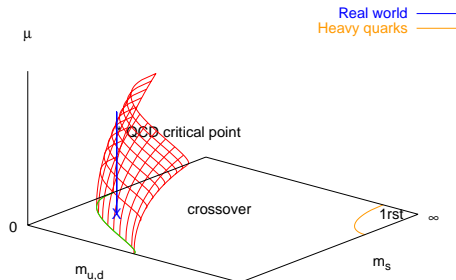


F&K keep $(m_q a)$ **fixed**, while $a(\beta_c)$ **increases with μ**
Lighter quarks at larger μ cause the phase transition
(dominant effect in our study)

- **Gavai & Gupta:** $\mu_E/T_E \lesssim 1$?
cutoff effects + **different theory** $N_f = 2$

Conclusions

- Need finer lattice $N_t \geq 6$ for **qualitative** determination of phase diagram



- $\frac{m_c(\mu)}{m_c(\mu=0)} = 1 + \mathcal{O}(1) \left(\frac{\mu}{\pi T}\right)^2 \rightarrow$ critical surface **almost vertical**
 - high sensitivity to quark masses (e.g. $m_u \neq m_d$)
 - $\mu_E/T_E \lesssim 1$ requires fine-tuning \rightarrow unnatural
- Moderate computer resources with **imaginary μ approach**