

***B* meson decays from moving-NRQCD on MILC fine lattices**

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Outline

1. $|V_{ub}| : B \rightarrow \pi l \nu$ semileptonic form factors $f_+(q^2)$, $f_0(q^2)$
2. lattice simulations with **NRQCD**
3. advantages of **moving-NRQCD**
4. determination of **renormalization constants** from high- β (weak coupling) simulations
5. preliminary results on B meson decays

Semileptonic decay $B \rightarrow \pi l \nu$

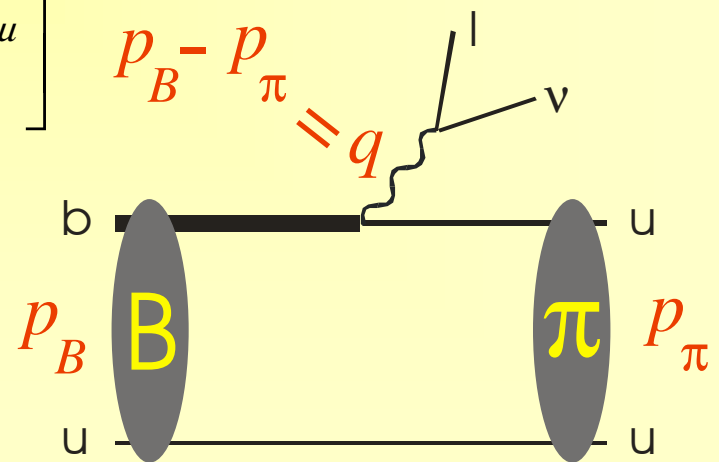
- CKM matrix elements ... new physics?
- $|V_{ub}|$ from semileptonic process $B \rightarrow \pi l \nu$

$$\frac{d\Gamma}{dq^2}(B \rightarrow \pi l \nu) = \frac{G_F^2}{24\pi^3} p_\pi^3 |V_{ub}|^2 |f_+(q^2)|^2$$

- **form factors $f_+(q^2)$, $f_0(q^2)$ from lattice QCD**

large $q^2 \rightarrow$ small recoil mom.
small $q^2 \rightarrow$ large recoil mom.

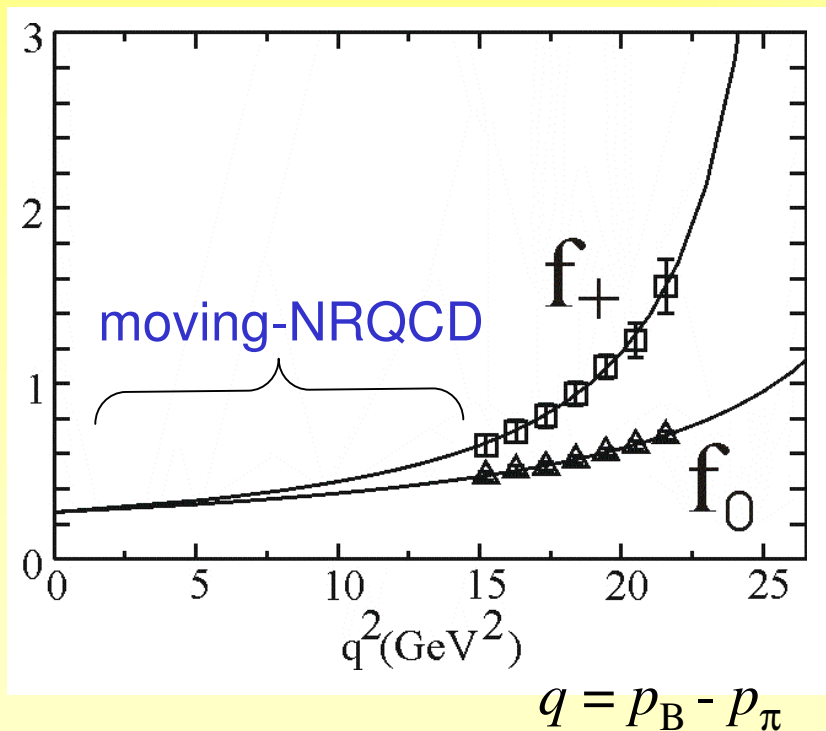
$$\langle \pi | V^\mu | B \rangle = f_+(q^2) \left[p_B^\mu + p_\pi^\mu - \frac{m_B^2 - m_\pi^2}{q^2} q^\mu \right] + f_0(q^2) \frac{m_B^2 - m_\pi^2}{q^2} q^\mu$$



$B \rightarrow \pi l \nu$ with NRQCD

➤ **NRQCD for the b quark**

➤ E. Gulez et al. HPQCD Collaboration. PRD73:074502 (2006)



➤ π can have very **large recoil momentum** (small q^2); but difficult to access this region in lattice calculations

1. discretization errors $\sim (ap)^2$
2. statistical errors grow exponentially as kinetic energy of the hadrons grows

Moving-NRQCD

- last transparency: ***B at rest*** + entire recoil momentum channeled into π
- ***solution***: B and π move in opposite direction
- ***problem***: large $(ap)^2$ errors from NRQCD
- ***moving-NRQCD***

[J.H. Sloan hep-lat/9710061; K.M. Foley & G.P.Lepage hep-lat/0209135]

$$p_b = m_b u + k ; \quad u = 4\text{-velocity of B} = p_B / m_B = \gamma(1, \vec{v})$$

treat the average momentum $m_b u$ exactly (like M_0 in NRQCD),
discretize the residual momentum $k \sim \text{small}$

Moving-NRQCD action

➤ action (boost the NRQCD action)

$$G_{x,t+1} = \left(1 - \frac{\delta H}{2}\right) \left(1 - \frac{H_0}{2n}\right)^n U_{x,\hat{t}}^\dagger \left(1 - \frac{H_0}{2n}\right)^n \left(1 - \frac{\delta H}{2}\right) G_{x,t}$$

$$H_0 = -iv \cdot \Delta^\pm - \frac{\Delta^{(2)}}{2M_0\gamma}$$

$$\begin{aligned} \delta H = & -\frac{1}{2m\gamma} \sigma \cdot B_m + \left(\frac{1}{2m\gamma} + \frac{1}{4n}\right) (v \cdot \Delta^\pm)^2 \\ & + \frac{1}{24m\gamma} \sum_i \Delta_i^{(4)} - \frac{1}{6m\gamma} (v \cdot \Delta^\pm) (\sum_i v_i \Delta_i^+ \Delta_i^\pm \Delta_i^-) \\ & + \frac{i}{6} \sum_i v_i \Delta_i^+ \Delta_i^\pm \Delta_i^- - \frac{i}{4mn\gamma} (v \cdot \Delta^\pm) \Delta^{(2)} \end{aligned}$$

**complete through
 $O(1/m)$**

$$B_m = \gamma \left(B - v \times E - \frac{\gamma}{\gamma+1} v(v \cdot B) \right)$$

Renormalization constants

- dispersion relation ***on the lattice***

$$E = E_0 + \sqrt{(Z_p p_0 + k)^2 + m_{kin}^2}$$

$$E_0 = c_1 \alpha_s + c_2 \alpha_s^2 + \dots$$

$$Z_p = 1 + c_1 \alpha_s + c_2 \alpha_s^2 + \dots$$

$$m_{kin} = Z_m m_0 = (1 + c_1 \alpha_s + c_2 \alpha_s^2 + \dots) m_0$$

- diagrammatic perturbation theory

1st order calculations finished

A. Dougall, C.T.H Davies, K.M. Foley & G.P.Lepage (lattice 2004)
also checked by R. Horgan *et al.* (University of Cambridge)

- ***numerical method (high- β simulations)***

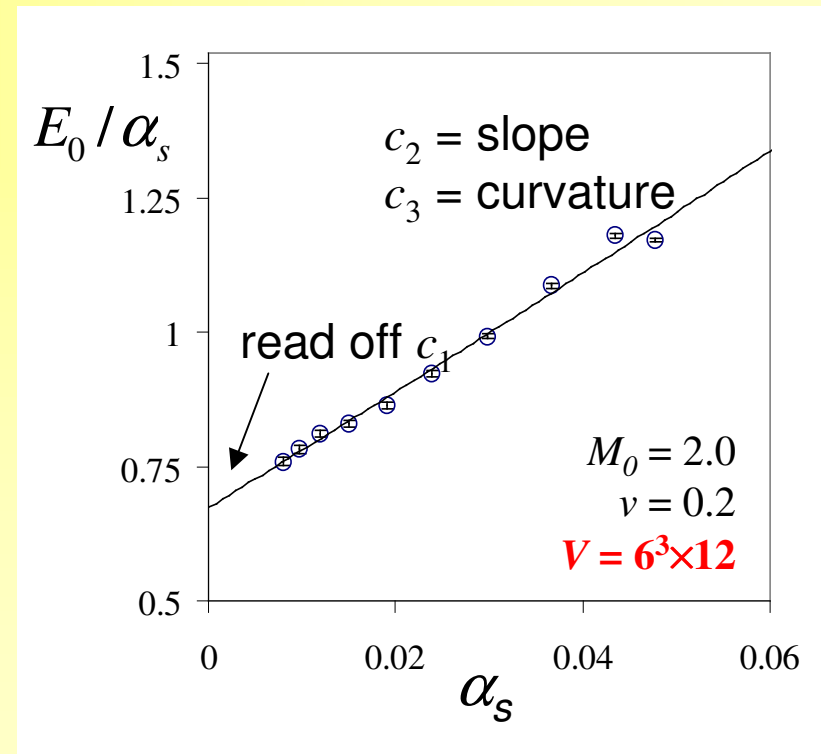
do several simulations at high- β (weak couplings) and fit the results to an expansion in α_s to extract the perturbative coefficients

High- β simulations

1. simulations at weak couplings
2. measure E_0 (from dispersion relation) in each case
3. fit the results to

$$E_0 = c_1 \alpha_s + c_2 \alpha_s^2 + \dots$$

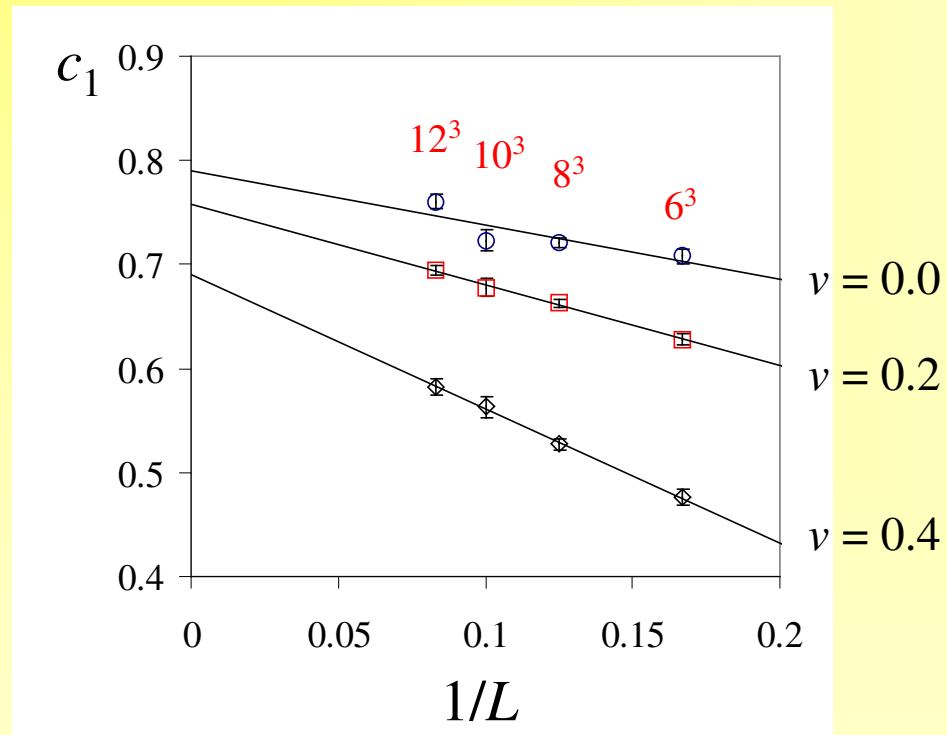
4. **read off PT coefficients**



quenched + H_0

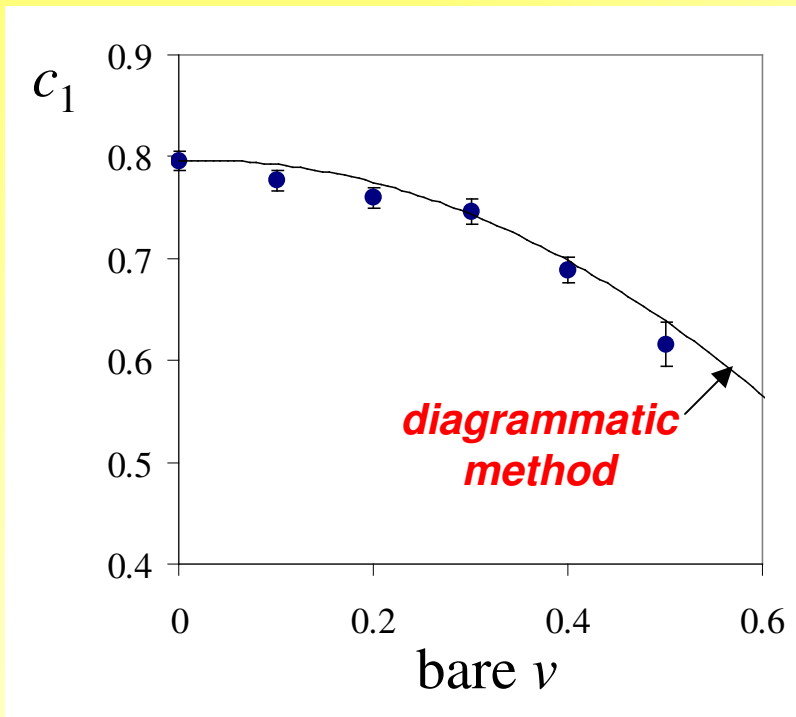
Difficulties

- difficulty #1 accurate measurements of observable
- difficulty #2 finite volume effects

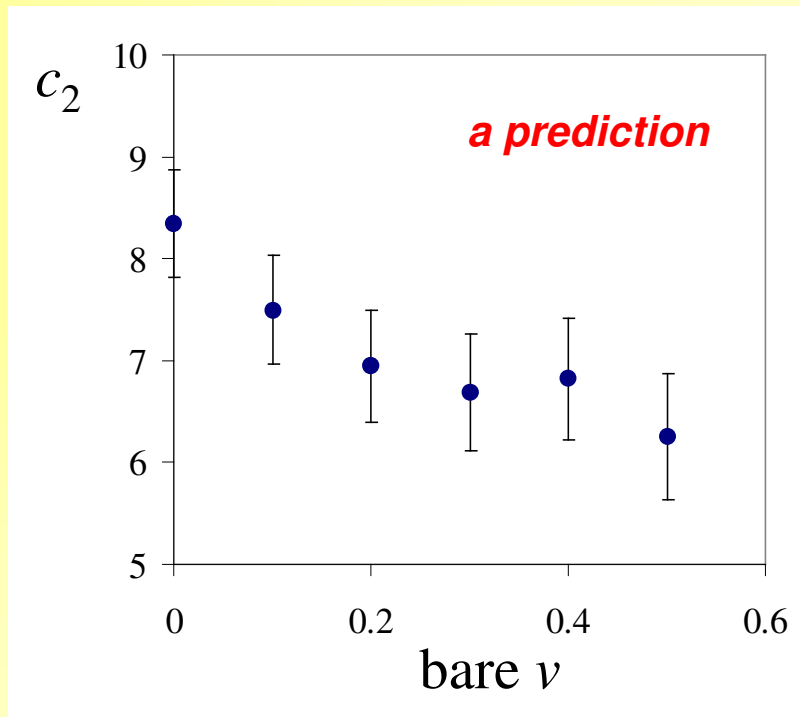


Results for E_0

1st order c_1

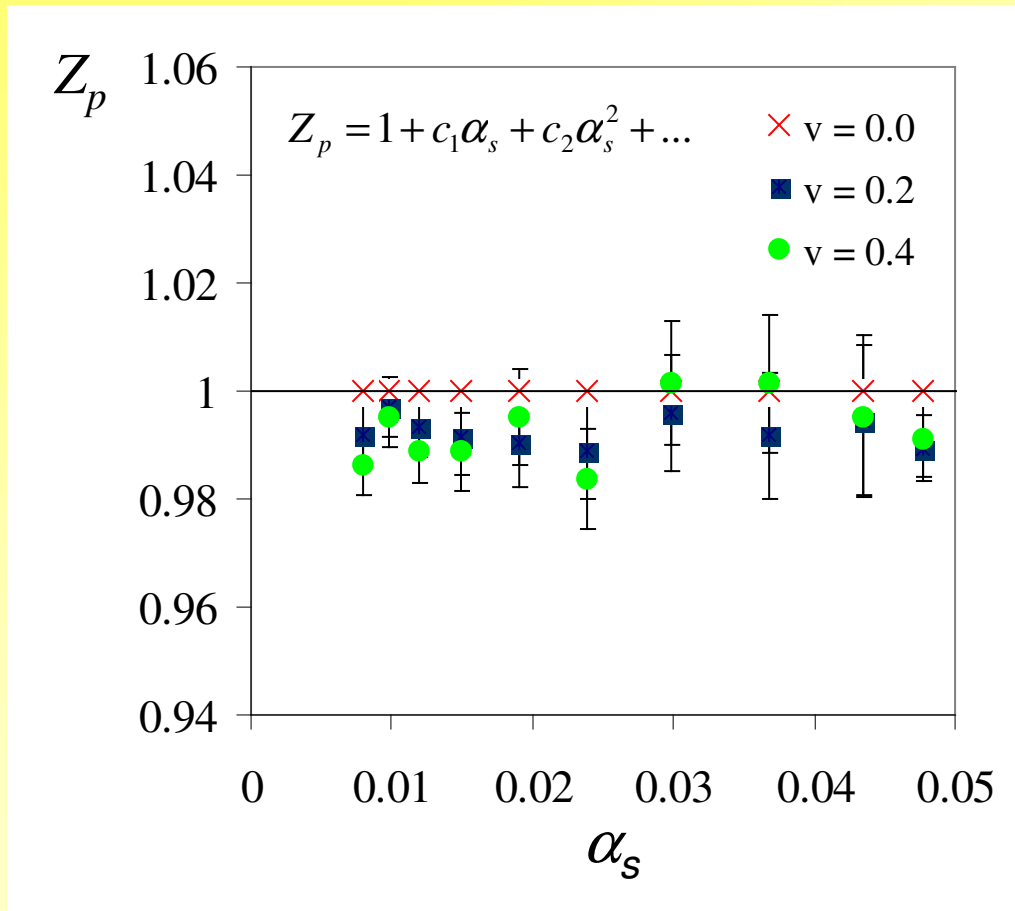


2nd order c_2



- high- β : ***much simpler*** compared to diagrammatic method ***at high orders***

Results for Z_p

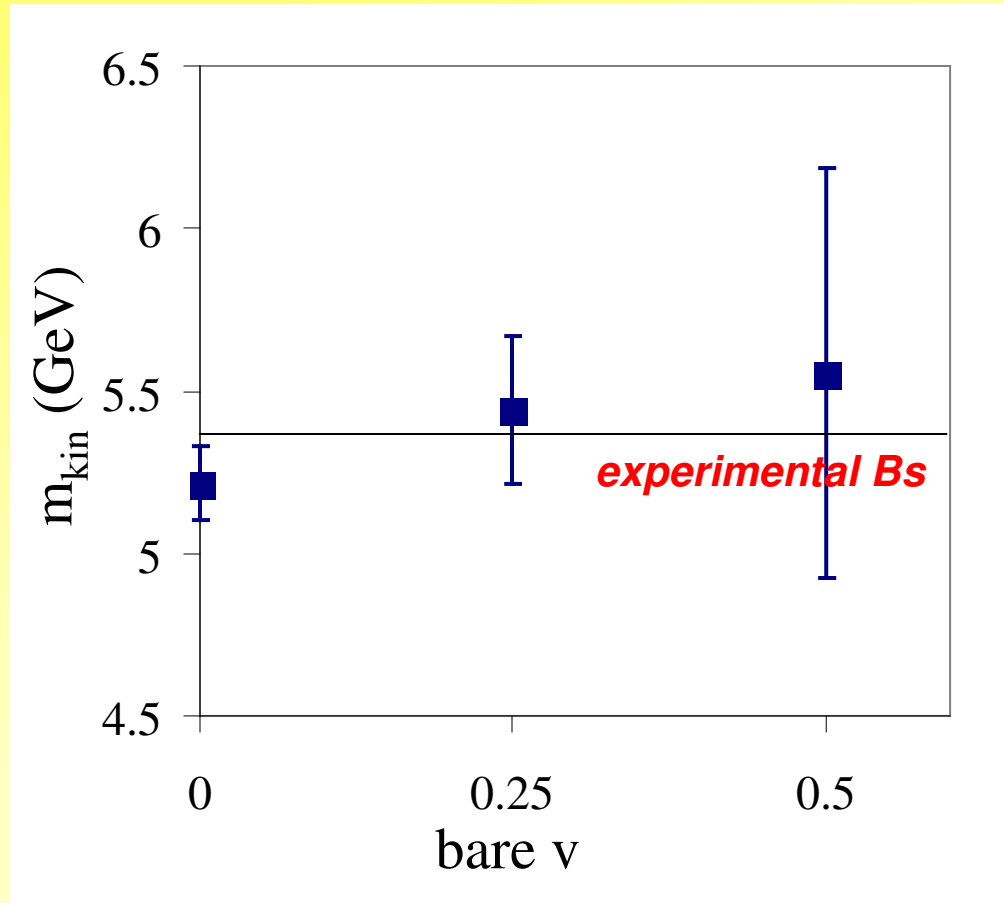


- error bars too large, difficult to do fits
- **good**: shows little renormalization, due to re-parameterization invariance

Simulations

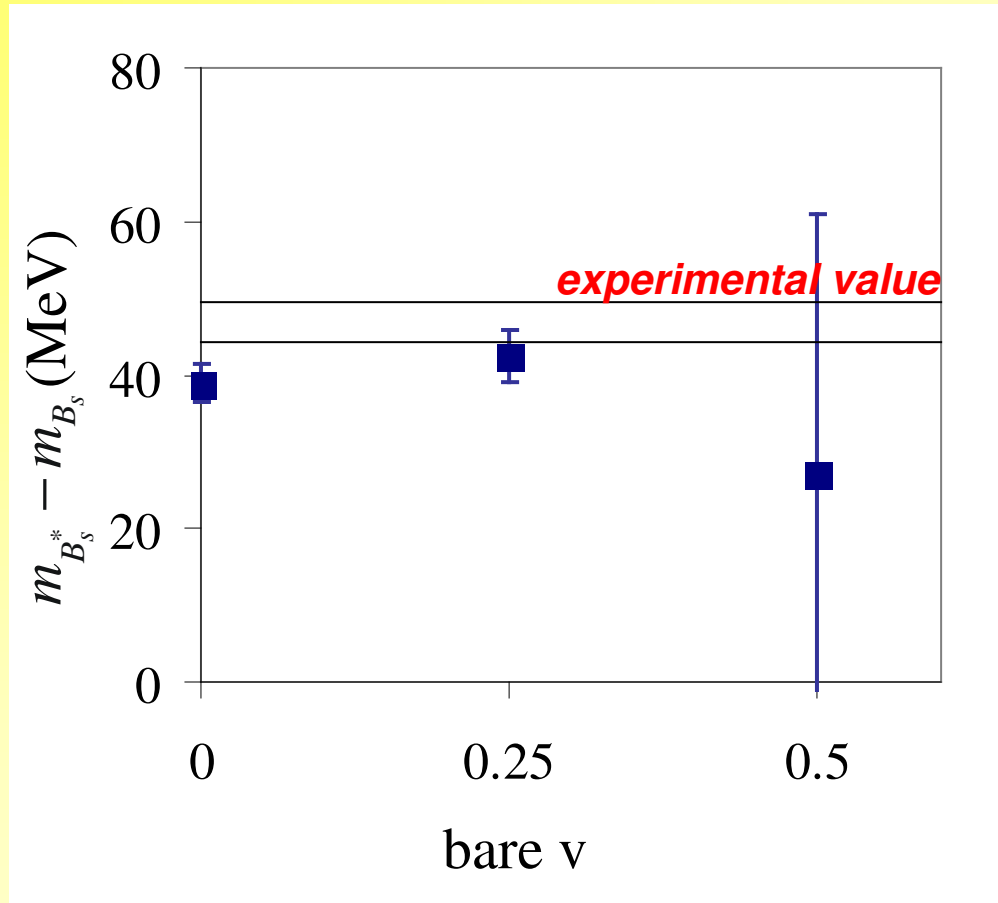
- decay of B meson (**2pt correlation function**)
- first study on **MILC fine lattices** ($a^{-1} \approx 2.258\text{GeV}$, $am_f = 0.0062$)
- asqtad light quark propagators ($am_q = 0.031 \sim \mathbf{Bs}$)
- $aM_0 = 1.95$ (see later)
- $v = (v_x, 0, 0)$; $v_x = 0, 0.25, 0.5$
- **we want to know**: how to set the parameters for the 3pt runs
 - dependence of m_{kin} on v ; need to tune M_0 ?
 - dependence of Z_p on v ; need to tune the bare v ?

m_{kin}

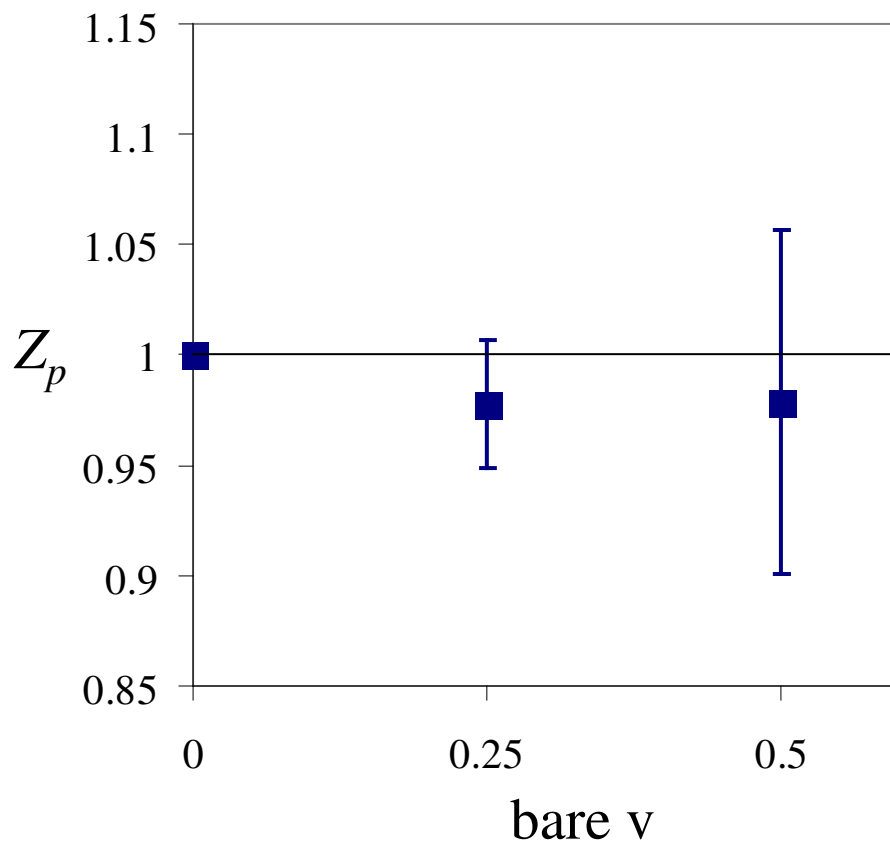


- input bare mass
 $aM_0 = 1.95$
- **good**: m_{kin} is quite stable as v increases
→ **no need to tune the bare mass**

$$m_{B_s^*} - m_{B_s}$$



Ext. Momentum Renormalization Z_p



- **good**: shows little renormalization (agree with PT)
→ **no need to tune the input velocity**
- small renormalization due to re-parameterization invariance

Conclusions

- **moving-NRQCD** → semileptonic form factors $f_+(q^2)$, $f_0(q^2)$ at **high recoil momentum (small q^2)**
- results for B decays on MILC fine lattices look good
- 3pt correlation function is in progress
- need renormalization constants: high- β simulations

**goal: high-precision determination
of $f_+(q^2)$, $f_0(q^2)$ and $|V_{ub}|$**