Quantum interference of large organic molecules

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The wave nature of matter is a key ingredient of quantum physics and yet it defies our classical intuition. First proposed by Louis de Broglie a century ago, it has since been confirmed with a variety of particles from electrons up to molecules. Here we demonstrate new high-contrast quantum experiments with large and massive tailor-made organic molecules in a near-field interferometer. Our experiments prove the quantum wave nature and delocalization of compounds composed of up to 430 atoms, with a maximal size of up to 60 Å, masses up to $m = 6,910$ AMU and de Broglie wavelengths down to $\lambda_{dB} = \frac{h}{mv} = 1$ pm. We show that even complex systems, with more than 1,000 internal degrees of freedom, can be prepared in quantum states that are sufficiently well isolated from their environment to avoid decoherence and to show almost perfect coherence.

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In many discussions on the foundations of physics, single-particle diffraction at a double slit or grating is regarded as a paradigmatic example for a highly non-classical feature of quantum mechanics, which has never been observed for objects of our macroscopic world. The quantum superposition principle has become of paramount importance also for the growing field of quantum information science. Correspondingly, research in many laboratories around the world is focusing on our understanding of the role of decoherence for increasingly complex quantum systems and possible practical or truly fundamental limits to the observation of quantum dynamics.

Here we report on a new leap in quantum interference with large organic molecules. In contrast to earlier successful experiments with internal molecular wave packets, our study focuses on the wave evolution in the centre of mass motion of the molecule as a whole, that is, pure de Broglie interference. We do this with compounds that have been customized to provide useful molecular beams at moderate temperatures. Figure 1 compares the size of two perfluoroalkylated nanospheres, PFNS8 and PFNS10, with a single C_{60} fullerene and it relates a single tetraphenylporphyrin molecule (TPP) to its complex derivatives TPPF84 and TPPF152. We demonstrate the wave nature of all these molecules in a three-grating near-field interferometer of the Kapitza-Dirac-Talbot-Lau type, as shown in Figure 2.

**Results**

**Experimental setup.** The particles are evaporated in a thermal source. Their velocity is selected using the gravitational free-fall through a sequence of three slits. The interferometer itself consists of three gratings G_1, G_2, and G_3 in a vacuum chamber at a pressure of p<10^{-4} mbar. The first grating is a SiN membrane with 90-nm wide slits arranged with a periodicity of d=266 nm. Each slit of G_1 imposes a constraint onto the transverse molecular position that, following Heisenberg’s uncertainty relation, leads to a momentum uncertainty. The latter turns into a growing delocalization and transverse coherence of the matter wave with increasing distance from G_1. The second grating, G_2, is a standing laser light wave with a wavelength of λ=532 nm. The interaction between the electric laser light field and the molecular optical polarizability creates a sinusoidal potential, which phase-modulates the incident matter waves. The distance between the first two gratings is chosen such that quantum interference leads to the formation of a periodic molecular density pattern 105 mm behind G_3. This molecular nanostructure is sampled by scanning a second SiN grating (G_3) identical to G_1 across the molecular beam while counting the number of the transmitted particles in a quadrupole mass spectrometer (QMS).

In extension to earlier experiments, we have added various technological refinements: the oven was adapted to liquid samples, a liquid-nitrogen-cooled chamber became essential to maintain the source pressure low, a new mass analyser allowed us to increase the detected molecular flux by a factor of four and many optimization cycles in the interferometer alignment were needed to meet all requirements for high-contrast experiments with very massive particles.

**Observed interferograms.** We recorded quantum interferograms for all molecules of Figure 1, as shown in Figure 3. In all cases the measured fringe visibility V, that is, the amplitude of the sinusoidal modulation normalized to the mean of the signal, exceeds the maximally expected classical moiré fringe contrast by a significant multiple of the experimental uncertainty. This is best shown for TPPF84 and PFNS8, which reached the highest observed interference contrast in our high-mass experiments so far, with individual scans up to V_{obs}=33% for TPPF84 (m=2,814 AMU) and V_{obs}=49% for PFNS8 at a mass of m=5,672 AMU. In addition, we have observed a maximum contrast of V_{obs}=17±4% for PFNS10 and V_{obs}=16±2% for TPPF152 (see Figure 3), in which our classical model predicts V_{class}=1%. This supports our claim of true quantum interference for all these complex molecules.

The most massive molecules are also the slowest and therefore the most sensitive ones to external perturbations. In our particle
of PFNS8, is characterized by a mean velocity of $v = 63 \text{ m s}^{-1}$ with a full width $\Delta v_{\text{FWHM}} = 13 \text{ m s}^{-1}$. The observed contrast of $V_{\text{isc}} = 49 \pm 6\%$ is in good agreement with the expected quantum contrast of $V_{\text{isc, quant}} = 51\%$ and is clearly discernible from the classically expected visibility of $V_{\text{isc}} < 1\%$. The stated uncertainty is the standard deviation of the fit to the data. (b) For PFNS10, the signal was too weak to allow a precise velocity measurement and quantum calculation. The oven position for these particles, however, limits the molecular velocity to $v < 80 \text{ m s}^{-1}$ and therefore allows us to define an upper bound to the classical visibility. (c) For TPPF84, we measure $v = 95 \text{ m s}^{-1}$ with $\Delta v_{\text{FWHM}} = 34 \text{ m s}^{-1}$. This results in $V_{\text{isc}} = 33 \pm 3\%$ with $V_{\text{isc, quant}} = 30\%$ and $V_{\text{isc}} < 1\%$. (d) The signal for TPPF152 is equally low compared with that of PFNS10. For this compound we find $V_{\text{isc}} = 16 \pm 2\%$, $V_{\text{isc, quant}} = 45\%$ and $V_{\text{isc}} < 1\%$.

Figure 3 | Quantum interferograms of tailor-made large organic molecules. Quantum interference well beyond the classical expectations has been observed for all molecules in the set. In all panels, the black circles represent the experimental result, the blue line is a sinusoidal fit to the shaded area which indicates the detector dark rate. (a) The beam of perfluoroalkylated nanospheres, PFNS8, is characterized by a mean velocity of $v = 63 \text{ m s}^{-1}$ with a full width $\Delta v_{\text{FWHM}} = 13 \text{ m s}^{-1}$. The observed contrast of $V_{\text{isc}} = 49 \pm 6\%$ is in good agreement with the expected quantum contrast of $V_{\text{isc, quant}} = 51\%$ and is clearly discernible from the classically expected visibility of $V_{\text{isc}} < 1\%$. The stated uncertainty is the standard deviation of the fit to the data. (b) For PFNS10, the signal was too weak to allow a precise velocity measurement and quantum calculation. The oven position for these particles, however, limits the molecular velocity to $v < 80 \text{ m s}^{-1}$ and therefore allows us to define an upper bound to the classical visibility. (c) For TPPF84, we measure $v = 95 \text{ m s}^{-1}$ with $\Delta v_{\text{FWHM}} = 34 \text{ m s}^{-1}$. This results in $V_{\text{isc}} = 33 \pm 3\%$ with $V_{\text{isc, quant}} = 30\%$ and $V_{\text{isc}} < 1\%$. (d) The signal for TPPF152 is equally low compared with that of PFNS10. For this compound we find $V_{\text{isc}} = 16 \pm 2\%$, $V_{\text{isc, quant}} = 45\%$ and $V_{\text{isc}} < 1\%$.

Comparison of theory and experiment. The experimental values have to be compared with the theoretical predictions based on a classical and a quantum model. The measured interference visibility is plotted as a function of the diffracting laser power $P$ in Figure 4 for TPPF84 (4a) and PFNS8 (4b). Our data are in very good agreement with the full quantum calculation and in clear discrepancy with the classical prediction. The abscissa scaling of the $V(P)$ curve is a good indicator for that. The quantum prediction mimics the classical curve qualitatively, but it is stretched in the laser power by a factor of about six (see Methods).

The laser power can be calibrated with an accuracy of $\pm 1\%$ but the abscissa also scales in proportion to the optical molecular polarizability and inversely with the vertical laser waist. The theoretical curves of Figure 4 are plotted assuming $\alpha_{\text{stat}} = 200 \text{ Å}^2 \times 4\pi\varepsilon_{\text{op}}$ for TPPF84 and $\alpha_{\text{stat}} = 190 \text{ Å}^2 \times 4\pi\varepsilon_{\text{op}}$ for PFNS8. These numbers have to be compared with the static polarizabilities computed using Gaussian09 (ref. 24). These are $\alpha_{\text{stat}} = 155 \text{ Å}^2 \times 4\pi\varepsilon_{\text{op}}$ for TPPF84 and $\alpha_{\text{stat}} = 200 \text{ Å}^2 \times 4\pi\varepsilon_{\text{op}}$ for PFNS8. A variation in the polarizability changes the horizontal scale of the plot as does a different laser waist. Both are bound by a relative uncertainty of less than 30%. A classical explanation is therefore safely excluded as an explanation for the experiments.

The quantitative agreement of the experimental and expected contrast is surprisingly good, given the high complexity of the particles. Various factors contribute to the remaining small discrepancies. The interference visibility is highly sensitive to apparatus vibrations, variations in the grating period on the level of 0.5 Å and a misalignment below 100 µrad in the grating roll angle.

Discussion

PFNS10 and TPPF152 contain 430 atoms covalently bound in one single particle. This is $\sim 1500$ atoms more than that in almost all previous experiments and it compares well with the number of atoms in small Bose–Einstein condensates (BEC), which, of course, operate in a vastly different parameter regime: The molecular de Broglie wavelength $\lambda_{\text{dB}}$ is about six orders of magnitude smaller than that of ultracold atoms and the internal molecular temperature exceeds typical BEC values ($T < 1 \mu K$) by about nine orders of magnitude. Although matter wave interference of BECs relies on the de Broglie wavelength of the individual atoms, our massive molecules always appear as single entities.

One can find various definitions in the literature for what exactly a Schrödinger cat should be and a number of intriguing experiments have reported the generation of photonic or atomic cat-states.
In as far as the term designates the quantum superposition of two macroscopically distinct states of a highly complex object, the molecules in our new experimental series are among the fattest Schrödinger cats realized to date. Schrödinger reasoned whether it is possible to bring a cat into a superposition state of being ‘dead’ and ‘alive’. In our experiment, the superposition consists of having all 430 atoms simultaneously ‘in the left arm’ and ‘in the right arm’ of our interferometer, that is, two possibilities that are macroscopically distinct. The path separation is about two orders of magnitude larger than the size of the molecules.

Schrödinger’s thought experiment originally also required the entanglement between a microscopic atom and the final state of the macroscopic cat. Such a mechanism is not needed to create the molecular superposition state in our experiment. Entanglement between a molecule and a microscopic probe particle does, however, occur in decoherence processes in which the quantum interaction with the environment reveals which-path information and destroys the interference pattern. Collisions with residual gas molecules, the emission of heat radiation and the absorption of blackbody radiation are among the most important decoherence mechanisms for interferometry with massive particles. We estimate that they lead to a visibility reduction of less than 1% under the conditions of the present experimental arrangement, in spite of the high internal molecular temperatures and substantial dipole fluctuations.

Specifically for PFN8, a宏观scopically realistic account of the decoherence processes predicts a visibility reduction of 10% only if the temperature of either the molecule or the radiation field exceeds 1,500 K, or if the residual nitrogen gas pressure exceeds 2×10⁻³ mbar.

In conclusion, our experiments reveal the quantum wave nature of tailor-made organic molecules in an unprecedented mass and size domain. They open a new window for quantum experiments with molecules in our new experimental series are among the fattest molecules in nature consisting of more than 1,000 internal degrees of freedom.

**Methods**

**Chemical synthesis.** The porphyrin derivatives were synthesized by the attachment of a highly fluoruous thiol to meso-tetra(pentafluorophenyl)porphyrin in a nucleaseogenic substitution reaction by applying a modified literature procedure. To assemble TPPF8, the commercially available 1H,1H,2H,2H-perfluorododecane-1-thiol as nucleaseogenic fluoruous part was introduced to the porphyrin unit. The branched thiou building block for TPPF152 was synthesized in three reaction steps. A reaction sequence including mono-functionalization of tris(bromomethyl)benzene with a protected thiol, introduction of two fluoruous ponytails and a final deprotection of the thiol functionality yielded the desired fluoruous thiou suitable for the envisaged substitution reaction. All target structures were purified by column chromatography and characterized by nuclear magnetic resonance spectroscopy and mass spectrometry (Supplementary Methods).

**Differences between the classical and quantum predictions.** The functional dependence of the interference fringe visibility on the laser power is qualitatively similar in both a classical and a full quantum treatment. As observed from the treatment described in ref. 23, the absissa scaling differs, however, by the factor \( \frac{\sin(\xi)}{\xi} \) where \( L_i \) is the distance between two consecutive gratings and \( L_i = dL_{\lambda, i} \) is the Talbot length. For the case of Figure 4, we find \( \frac{\sin(\xi)}{\xi} \approx 5.9 \). The experimental data are in clear agreement with the quantum model.

**Equipment.** The deflecting laser beam is generated by a Coherent Verdi V18 laser at 532 nm. The QMS is an Elsle CMS with a rod diameter of 9.5 mm, operated at a radio frequency of 440 kHz. The SiN gratings in G₁ and G₂ were made by Dr Tim Savas, nm² LLC & MIT Cambridge.

**References**


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Author contributions
S.G. and S.E. performed all interference experiments as well as the analysis of the data with important contributions by M.T. M.A. contributed at various stages of the experiment. J.T. synthesized, purified and analysed the porphyrin derivatives based on a design developed together with M.M. P.F. provided the perfluoroalkylated nanospheres. M.A. and M.M. initiated and coordinated the experiments. S.N. and K.H. participated in the interpretation of the data. M.A., S.G. and S.N. wrote the paper. All authors discussed the results and commented on the manuscript.

Additional information
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